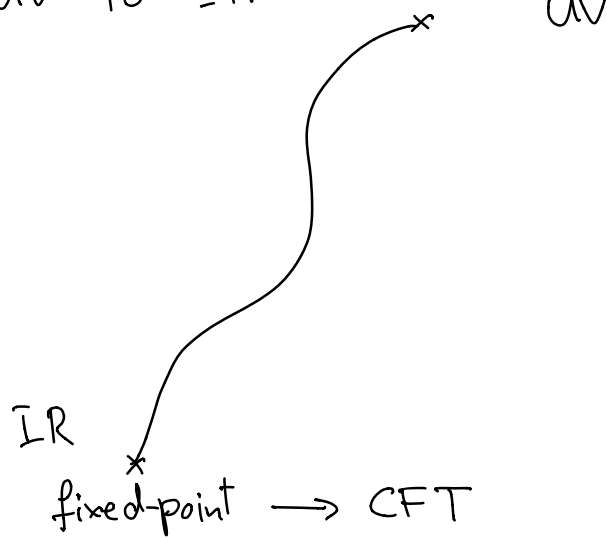
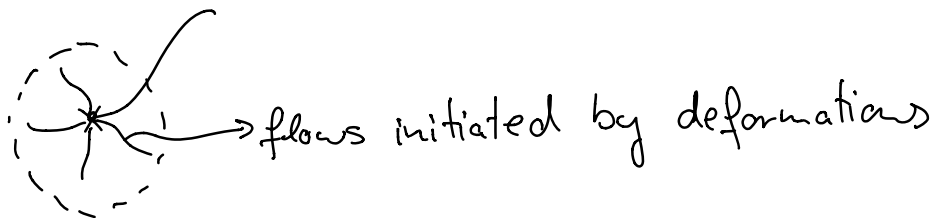


§2. Deformations of Superconformal Theories

QFT's are renormalization group (RG) flows
from UV to IR



Given a fixed point, we want to analyze its deformations



3 classes:

1) Adding local operators to the Lagrangian

$$\delta\mathcal{L} = g O$$

running coupling constant

local operator
defined at $g=0$

2) Gauging a global symmetry:
 gauge a continuous flavor symmetry
 with conserved current j_μ
 \rightarrow may be obstructed by anomalies

3) Moving onto a moduli space of vacua:
 In $d > 2$ spacetime dimensions CFT may
 possess non-trivial moduli space of vacua.
 \rightarrow breaks conformal symmetry
 spontaneously (restored at origin)

We will focus on deformation 1)

\mathcal{O} must reside in representations of $SO(d, 2)$
 \rightarrow labelled by weights under $SO(d) \times SO(2)$

Have unique operator \mathcal{O} of lowest dimension Δ_0
 known as conformal primary

$$\begin{array}{ccc}
 K_\mu \mathcal{O} = 0 & & P_\mu \mathcal{O} \\
 \uparrow & & \uparrow \\
 \text{of scaling dimension } -1 & & \text{descendants of} \\
 & & \text{scaling dimension } \Delta_0 + 1 \\
 & & (P_\mu \sim \partial_\mu)
 \end{array}$$

\mathcal{O} cannot be written as $\partial_\mu \mathcal{O}'$

Let us denote Δ_0 and $SO(d)$ weights of \mathcal{O} by
 $L_{\mathcal{O}}$. There is natural inner product $\langle \mathcal{O}' | \mathcal{O}'' \rangle$

require $\langle \mathcal{O}' | \mathcal{O}'' \rangle \geq 0$

→ unitarity bound $\Delta_{\mathcal{O}} \geq f(L_{\mathcal{O}})$

→ null-states when bound is saturated

A deformation $\delta \mathcal{L} \sim g \mathcal{O}$ must be

- a conformal primary
(descendants $\partial_n \mathcal{O}$ lead to total derivatives)
- $SO(d)$ scalar → preserves Lorentz symmetry

→ $\langle \square \mathcal{O} | \square \mathcal{O} \rangle \geq 0 \Rightarrow \Delta_{\mathcal{O}} \geq \frac{d-2}{2}$

(recall $\Delta_{\mathcal{O}} = h_1 + \frac{d-2}{2}$)

↑
highest weight of $S(d)$
 $S(d-2)$ -weights free

bound is saturated for $\square \mathcal{O} = 0$

Deformation types:

- "relevant deformations" ($\Delta_{\mathcal{O}} < d$):

CFT at $g=0$ UV fixed point of RG-flow

→ g grows in the IR (perturbation theory breaks down)

- "Irrelevant deformations" ($\Delta_{\mathcal{O}} > d$):

CFT is at IR fixed-point of RG-flow

→ g flows to zero

- "Marginal deformations" ($\Delta_{\mathcal{O}} = d$):

preserve conformal invariance → lead to nearby fixed point

superconformal theories:

Want to look at deformations preserving
Poincaré \mathcal{Q} -supersymmetries
(not necessarily \mathcal{S} -symmetries)

Recall that we have the following SCA's:

$$d=3 \quad \text{osp}(\mathcal{N}|4) \supset \text{SO}(3,2) \times \text{SO}(\mathcal{N})_R$$

$$d=4 \quad \begin{cases} \text{sl}(2,2|\mathcal{N}) \supset \text{SO}(4,2) \times \text{SU}(\mathcal{N})_R \times \text{U}(1)_R, \mathcal{N} \neq 4 \\ \text{psl}(2,2|4) \supset \text{SO}(4,2) \times \text{SU}(4)_R, \mathcal{N}=4 \end{cases}$$

$$d=5 \quad \text{F}(4) \supset \text{SO}(5,2) \times \text{SU}(2)_R, \mathcal{N}=1$$

$$d=6 \quad \text{osp}(6,2|\mathcal{N}) \supset \text{so}(6,2) \times \text{sp}(2\mathcal{N})_R$$

\mathcal{N} denotes number of d -supercharges

$N_{\mathcal{Q}}$ denotes total number of supercharges

In $d=3,4,5,6$ minimal $\mathcal{N}=1$ supersymmetry
corresponds to $N_{\mathcal{Q}} = 2, 4, 8, 8$

Notation: Will consider conformal primaries

labelled by $L_{\mathcal{O}}, R_{\mathcal{O}}, \Delta_{\mathcal{O}}$

\nearrow Lorentz-charge \uparrow R-sym charge \nwarrow scaling dim

$$\mathcal{O} \in [L_{\mathcal{O}}]_{\Delta_{\mathcal{O}}}^{(R_{\mathcal{O}})}$$

superconformal primaries:

a conformal primary \mathcal{V} with

- lowest scaling dimension $\Delta_{\mathcal{V}}$
- irreducible under R -sym

→ annihilated by S (scaling dim $-\frac{1}{2}$)
and K_n (scaling dim. -1)

other conformal primaries are superconformal descendants of \mathcal{V} (by action of Q)

Unitarity bound:

$$\Delta_{\mathcal{V}} \geq f(L_n, R_n)$$

saturation leads to null-states

→ descendants form sub-representation
notation: "short" representation

Deformations:

conformal primary \mathcal{O} with

$$Q\mathcal{O} = \mathcal{D}_n(\dots)$$

→ must transform to descendant
under action of Q

→ cannot be superconformal primary

→ must be a "top" component

$Q^l \mathcal{V}$ are "level" l conformal primaries
(l nested anti-commutators)

$\{Q_i, Q_j\} \sim 0, \quad i, j = 1, \dots, N_Q$
 (we are dropping ∂_{μ} -terms, correspond to conformal descendants)

We have:

$0 \leq l \leq l_{\max}$ with $l_{\max} \leq N_Q$

→ saturated for "long" multiplets (no "null"-states)

$Q^{N_Q} \mathcal{V}$ is top component

Lorentz- and R-sym singlet

"D-term" deformation $\mathcal{L}_D = Q^{N_Q} \mathcal{V}$

$\Delta(\mathcal{L}_D) = \frac{1}{2} N_Q + \Delta_{\mathcal{V}}$

example: Kähler potential $\int d^4\theta K(\Phi, \bar{\Phi})$ in $d=4$ ($\mathcal{N}=1$)

→ F-terms correspond to "short" multiplets (BPS)

\mathcal{V}_{BPS} is annihilated by half of supercharges

$\mathcal{L}_F = Q^{\frac{1}{2} N_Q} \mathcal{V}_{\text{BPS}}$

example: chiral superpotential $\int d^2\theta W$

→ sporadic very short multiplets

example: consider the stress-tensor multiplet in

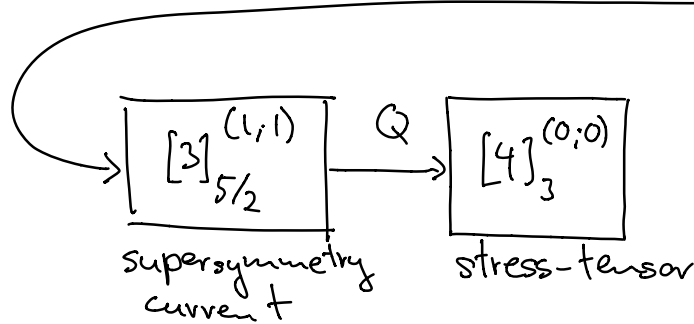
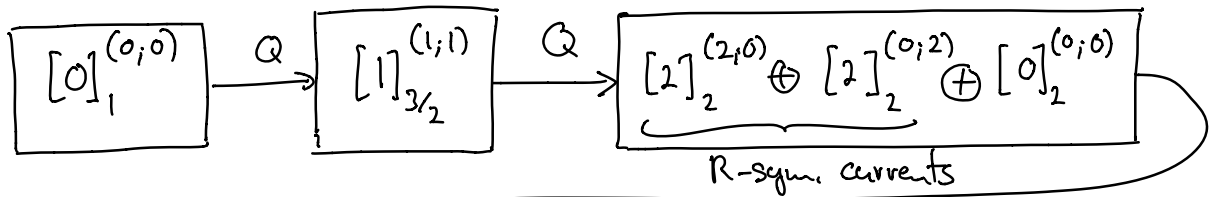
$\mathcal{N}=4$ $d=3$ sEFT's

Lorentz-sym.: $SU(2)$

R-sym.: $SU(2)_R \times SU(2)'_R$

$\rightarrow [j]_{\Delta}^{(R; R')}$

$Q_{\alpha}^{i,i'}$ transform as $[1]_{1/2}^{(1;1)} \rightarrow$ conformal prim. decomp.:



$[0]_2^{(0;0)}$ at level 2 is Lorentz-scalar

$$Q [0]_2^{(0;0)} \sim 0$$

$\rightarrow [0]_2^{(0;0)}$ gives relevant deformation with scaling dimension $\Delta = 2$.

Table of deformations in various dimensions:

see next page

d	\mathcal{N}	Relevant	Marginal	Irrelevant Δ_{\min}
3	1	D-term	D-term	$\Delta_{\min} > 3$
	2	Flavor current, F.t.	F-term	$\Delta_{\min} > 3$
	3	Flavor current	—	4
	4	stress t., fl. curr.	—	4
	5, 6	stress tensor	—	5
	8	stress tensor	—	6
4	1	F-term	F-term	$\Delta_{\min} > 4$
	2	fl. curr., F-term	F-term	$\Delta_{\min} > 4$
	3	—	—	$\Delta_{\min} > 4$
	4	—	stress tensor	8
5	1	flavor curr.	—	8
6	(1,0)	—	—	10
	(2,0)	—	—	12